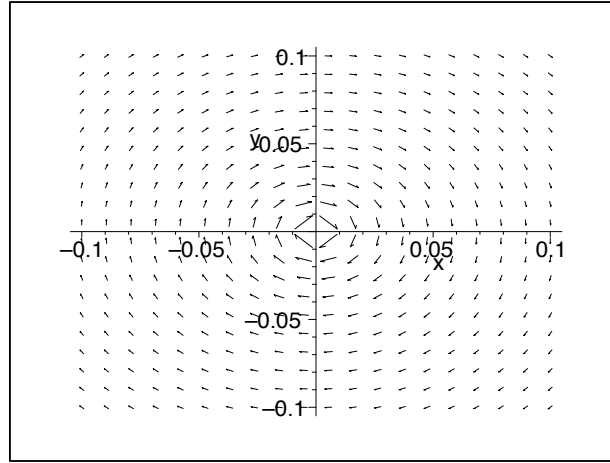


1. The volume can be calculated using the scalar triple product, using the three vectors emanating from any vertex of the parallelepiped. For example, [3]

$$\begin{aligned} V &= (\mathbf{v} - \mathbf{a}) \cdot (\mathbf{b} - \mathbf{a}) \times (\mathbf{d} - \mathbf{a}) = \mathbf{v} \cdot \mathbf{b} \times \mathbf{d} \quad (\text{since } \mathbf{a} = \mathbf{0}) \\ &= (1, -1, 1) \cdot (0, 0, 3) = 3. \end{aligned}$$

2. (a) $\mathbf{V}(\mathbf{r}) = \frac{(y\hat{\mathbf{i}} - x\hat{\mathbf{j}})e^{-a(x^2+y^2)^{3/2}}}{x^2 + y^2} = \frac{\rho \sin \phi \hat{\mathbf{i}} - \rho \cos \phi \hat{\mathbf{j}}}{\rho^2} e^{-a\rho^3} = -\frac{\hat{\phi}}{\rho} e^{-a\rho^3}. \quad [3]$

- (b) To plot the vector *field* it is necessary to show how the vector field varies in both direction and magnitude, from point to point. This means drawing vectors all



over the place!

[2]

- (c) In polars: to calculate $\nabla \times \mathbf{V}$, note that there is only a ϕ component to \mathbf{V} ; furthermore, this component only depends on ρ . Hence, we only need [looking up the identity for curl and extracting the relevant component]

$$\nabla \times \mathbf{V} = \hat{\mathbf{k}} \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho V_\phi) = \hat{\mathbf{k}} \frac{1}{\rho} \frac{\partial}{\partial \rho} (-e^{-a\rho^3}) = 3a\rho e^{-a\rho^3} \hat{\mathbf{k}}.$$

In cartesians it's a slog. There's only a z component (To see this, note that the x and y components of \mathbf{V} do not depend on z and that there is no z component of \mathbf{V} ; and then stare at the usual determinant). The chain rule has to be applied many times. [4]

$$\begin{aligned} (\nabla \times \mathbf{V})_z &= \frac{\partial V_x}{\partial y} - \frac{\partial V_y}{\partial x} = \left\{ \left[\frac{1}{x^2 + y^2} - \frac{y(2y)}{(x^2 + y^2)^2} + \frac{y}{x^2 + y^2} (3a/2)(2y)(x^2 + y^2)^{1/2} \right] \right. \\ &\quad \left. - (-) \left[\frac{1}{x^2 + y^2} - \frac{x(2x)}{(x^2 + y^2)^2} + \frac{x}{x^2 + y^2} (3a/2)(2x)(x^2 + y^2)^{1/2} \right] \right\} e^{-a(x^2 + y^2)^{3/2}} \\ &= \left[\frac{2(x^2 + y^2) - 2x^2 - 2y^2}{(x^2 + y^2)^2} + \frac{3a(x^2 + y^2)(x^2 + y^2)^{1/2}}{x^2 + y^2} \right] e^{-a(x^2 + y^2)^{3/2}} \\ &= 3a\sqrt{x^2 + y^2} e^{-a(x^2 + y^2)^{3/2}}. \end{aligned}$$

This is the same as the result in polar coordinates, since $\rho = \sqrt{x^2 + y^2}$.

3. Given $h(x, y) = (x^2 + 1)/(1 + x^2 + y^3)$.

(a) [3]

$$\begin{aligned}\nabla h &= \left(\hat{\mathbf{i}} \frac{\partial}{\partial x} + \hat{\mathbf{j}} \frac{\partial}{\partial y} \right) h = \hat{\mathbf{i}} \frac{2x(1 + x^2 + y^3) - (1 + x^2)(2x)}{(1 + x^2 + y^3)^2} + \hat{\mathbf{j}} \frac{-3y^2(1 + x^2)}{(1 + x^2 + y^3)^2} \\ &= \frac{2xy^3\hat{\mathbf{i}} - 3y^2(1 + x^2)\hat{\mathbf{j}}}{(1 + x^2 + y^3)^2}.\end{aligned}$$

(b) To find the slope in the direction of $-\hat{\mathbf{i}} - \hat{\mathbf{j}}$, we need to project the gradient ∇h along a unit vector in this direction. Hence, [2]

$$\begin{aligned}\text{slope} &= -\frac{1}{\sqrt{2}}(\hat{\mathbf{i}} + \hat{\mathbf{j}}) \cdot \nabla h \Big|_{x=1, y=2} = -\frac{1}{\sqrt{2}} \frac{2xy^3 - 3y^2(1 + x^2)}{(1 + x^2 + y^3)^2} \Big|_{x=1, y=2} \\ &= \frac{1}{\sqrt{2}} \frac{-16 + 24}{(1 + 1 + 8)^2} = \frac{1}{\sqrt{2}} \frac{8}{100} = \frac{2}{25\sqrt{2}}.\end{aligned}$$

(c) The maximum slope at this point is the magnitude of the gradient at this point, or [2]

$$\sqrt{\nabla h \cdot \nabla h} = \left| \frac{16}{100}\hat{\mathbf{i}} - \frac{24}{100}\hat{\mathbf{j}} \right| = \frac{\sqrt{16^2 + 24^2}}{100} = \frac{2\sqrt{13}}{25}.$$

4. The Laplacian operator, in cartesian coordinates, is $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$. So: [4]

$$(a) \nabla^2(x^2 + 2xy + 3ze^{-(x+y)}) = [2 + 0 + 0] + [0 + 0 + 0] + [3ze^{-(x+y)} + 3ze^{-(x+y)} + 0] = 2 + 6ze^{-(x+y)}.$$

$$(b) \nabla^2[e^{-5x} \sin 4y \cos 3z] = (25 - 16 - 9)e^{-5x} \sin 4y \cos 3z = 0.$$

5. Maxwell's Equations for electric and magnetic fields \mathbf{E} and \mathbf{B} in free space are

$$\nabla \cdot \mathbf{B} = 0 \quad (1)$$

$$\nabla \cdot \mathbf{E} = 0 \quad (2)$$

$$\nabla \times \mathbf{E} + \frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = 0 \quad (3)$$

$$\nabla \times \mathbf{B} - \frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} = 0 \quad (4)$$

(a) Take the curl of Equation ??:

$$\begin{aligned}\nabla \times \nabla \times \mathbf{E} + \frac{1}{c} \nabla \times \frac{\partial \mathbf{B}}{\partial t} &= \underbrace{\nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}}_{\text{identity}} + \underbrace{\frac{1}{c} \frac{\partial}{\partial t} \nabla \times \mathbf{B}}_{\text{change derivative order}} = 0 \\ \underbrace{0 - \nabla^2 \mathbf{E}}_{\text{use } \nabla \cdot \mathbf{E} = 0} + \frac{1}{c} \frac{\partial}{\partial t} \underbrace{\frac{1}{c} \frac{\partial \mathbf{E}}{\partial t}}_{\text{Use Eq. ??}} &= 0 \implies c^2 \nabla^2 \mathbf{E} = \frac{\partial^2}{\partial t^2} \mathbf{E}.\end{aligned}$$

This is the full wave equation in three dimensional space (rather than the one dimensional cousin for a wave on a string, $\frac{\partial^2 h}{\partial t^2} = c^2 \frac{\partial^2 h}{\partial x^2}$). Note that this is actually *three* wave equations, one for each component of \mathbf{E} ! A similar calculation can be performed for \mathbf{B} . [3]

(b) If $\mathbf{B} = \mathbf{B}_0 \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)$ and $\mathbf{E} = \mathbf{E}_0 \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)$, then Equation ?? implies:[3]

$$\begin{aligned} \nabla \times [\mathbf{E}_0 \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)] + \frac{1}{c} \frac{\partial [\mathbf{B}_0 \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)]}{\partial t} &= 0 \\ \underbrace{[\nabla \sin(\mathbf{k} \cdot \mathbf{r} - \omega t)] \times \mathbf{E}_0}_{\mathbf{E}_0 \text{ constant}} + \frac{1}{c} \mathbf{B}_0 \frac{\partial}{\partial t} [\sin(\mathbf{k} \cdot \mathbf{r} - \omega t)] &= 0 \\ \left[\mathbf{k} \times \mathbf{E}_0 + \frac{1}{c} (-\omega) \mathbf{B}_0 \right] \cos(\mathbf{k} \cdot \mathbf{r} - \omega t) &= 0 \end{aligned}$$

This must be true for all points \mathbf{r} and all times t ; hence we must have $\mathbf{k} \times \mathbf{E}_0 = \omega \mathbf{B}_0 / c$. That is, the magnetic and electric fields are polarized perpendicular to each other and, since we know that $\mathbf{k} \perp \mathbf{E}_0$ from Maxwell Equation (??), and $k = \omega / c$ [this can be found from the wave equation derived in the previous part], the electric and magnetic fields are equal in magnitude, in these units.